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¹H and ¹⁹F NMR Study of Ammonium Ion Motion in Ammonium Trifluorostannate (II)

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The temperature dependences of the second moment and spin-lattice relaxation time of $^1{\rm H}$ and $^{19}{\rm F}$ NMR were measured on ammonium trifluorostannate (II) NH₄SnF₃. It was found that the NH₄ ions having a C₃ symmetry in the crystal undergo overall reorientations with an activation energy of 22 kJ mol $^{-1}$.

Ammonium trifluorostannate (II) $\mathrm{NH_4SnF_3}$ forms trigonal crystals with the space group R $\overline{3}$ at room temperature [1, 2]. Both the nitrogen and tin atoms in isolated $\mathrm{NH_4^+}$ and $\mathrm{SnF_3^-}$ ions, respectively, are on the $\mathrm{C_3}$ symmetry axis of the crystal. One of the N – H bonds in an $\mathrm{NH_4^+}$ ion is on the $\mathrm{C_3}$ axis and directs to the center of a $\mathrm{SnF_3^-}$ trigonal pyramid, making a trifurcated hydrogen bond. The N ... F distances are 300 pm. The remaining three hydrogens of the $\mathrm{NH_4^+}$ ion are normally hydrogen-bonding with fluoride ions belonging to different $\mathrm{SnF_3^-}$ ions. The N ... F distances are 275 pm. Thus, the $\mathrm{NH_4^+}$ ion in $\mathrm{NH_4SnF_3}$ forms two kinds of very different hydrogen bonds. In order to study the motion of the $\mathrm{NH_4^+}$ ions forming such unique hydrogen bonds, we measured the second moments M_2 and spin-lattice relaxation times T_1 of $^1\mathrm{H}$ and $^{19}\mathrm{F}$ NMR of this complex.

Ammonium trifluorostannate (II) was crystallyzed from an aqueous solution of SnF₂ and a slight excess of NH₄HF₂ [2, 3]. The colorless crystals were filtered off, washed with methanol and dried over NaOH in a desiccator. NH₄SnF₃ thus obtained was identified by powder X-ray diffraction.

thus obtained was identified by powder X-ray diffraction. Continuous-wave ^1H and ^{19}F NMR spectra were recorded on a JEOL JNM-MW 40S spectrometer. ^1H and ^{19}F spin-lattice relaxation times T_1 were measured at 20 and 18.814 MHz (the external magnetic field is ca. 0.47 T for both nuclei), respectively, by using pulsed NMR spectrometers already described [4, 5]. The usual 180° -t- 90° pulse sequence was used for the determination of T_1 . Temperatures were determined by a copper-Constantan thermocouple within an accuracy of \pm 1 K for the T_1 measurements and of \pm 2 K for the M_2 measurements.

The second moments of ^{1}H and ^{19}F NMR absorptions ($M_{2\text{H}}$ and $M_{2\text{F}}$, respectively) were measured from 77 K to room temperature. At 77 K, $M_{2\text{H}}$ and $M_{2\text{F}}$ were 50×10^{-8} and 11×10^{-8} T², respectively. With increasing temperature, each M_{2} decreased in a temperature range 120-180 K, and reached a plateau value above 180 K. The plateau values of $M_{2\text{H}}$ and $M_{2\text{F}}$ were 4×10^{-8} and 7×10^{-8} T², respectively. The decrease of $M_{2\text{H}}$ is very large compared with that of

The decrease of $M_{2\rm H}$ is very large compared with that of $M_{2\rm F}$, suggesting the occurrence of isotropic reorientation of the NH₄⁺ ions. Therefore, the theoretical $M_{2\rm H}$ and $M_{2\rm F}$ were calculated for the models of rigid lattice and of isotropically

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reorienting NH $_{4}^{+}$ ions. Magnetic dipolar interactions between 1 H and 1 H, 1 H and 19 F, and 19 F and 19 F nuclei were taken into account in the M_{2} calculations. All interproton distances in an NH $_{4}^{+}$ ion were assumed to be 170 pm for the rigid lattice model, and for the reorienting NH $_{4}^{+}$ ion model the interionic contribution to M_{2} was estimated by placing four NH $_{4}^{+}$ ion protons at the nitrogen site. The results were $M_{2\rm H}=53.1\times10^{-8}$ and $M_{2\rm F}=12.5\times10^{-8}$ T $_{2}^{-}$ for the rigid lattice model and $M_{2\rm H}=2.7\times10^{-8}$ and $M_{2\rm F}=6.2\times10^{-8}$ T $_{2}^{-}$ for the reorienting NH $_{4}^{+}$ ion model. These calculated M_{2} values are in good agreement with the corresponding values observed at low and high temperatures, indicating that the NH $_{4}^{+}$ ions are rigidly fixed below 120 K and undergo isotropic reorientation above 180 K.

Figure 1 shows the temperature dependences of ^{1}H and ^{19}F T_{1} (T_{1H} and T_{1F} , respectively). A T_{1H} minimum of 4 ms and a T_{1F} minimum of 28 ms were located at 210 and 195 K, respectively. Below ca. 190 K, nonexponential decays of the magnetization recovery for the T_{1} measurements were observed for both nuclei. This nonexponentiality is assignable to the heteronuclear dipolar cross coupling between ^{1}H and ^{19}F nuclei [6–8]. In this case, the magnetization recovery curve was assumed to be the superimposition of two exponential decays, and a short (T_{1s}) and a long (T_{11}) component were manually determined as shown in Figure 1. Experimentally equal values of each component were obtained for both

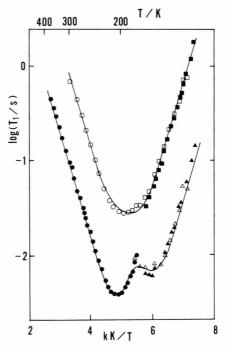


Fig. 1. Temperature dependences of spin-lattice relaxation times T_1 of ${}^1\mathrm{H}$ (\bullet , \blacksquare , and \blacktriangle) and ${}^{19}\mathrm{F}$ (\bigcirc , \square , and \triangle) NMR measured at 20 and 18.814 MHz, respectively, for NH₄SnF₃. Below ca. 190 K, nonexponential recovery of magnetization in the T_1 measurements was observed for both nuclei, yielding a long (\blacksquare and \square) and a short (\blacktriangle and \triangle) T_1 component.

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nuclei. A T_{1s} minimum of ca. 7 ms was observed near 165 K.

By referring the results of M_2 , the observed T_1 was analyzed by the model of the NH₄ ion reorientation. The relaxation rates (T_1^{-1}) of interacting ¹H and ¹⁹F nuclei in NH₄SnF₃ are the eigen values of the relaxation rate matrix [6–8]

$$\begin{pmatrix} R_{\rm HH} & R_{\rm HF} \\ R_{\rm FH} & R_{\rm FF} \end{pmatrix}. \tag{1}$$

The components of the relaxation matrix are given as

$$\begin{split} R_{\rm HH} &= (2/3) \, \gamma_{\rm H}^2 \, \Delta M_2 \, ({\rm H} - {\rm H}) \, f \, (\omega_{\rm H}, \tau) \\ &+ (1/2) \, \gamma_{\rm H}^2 \, \Delta M_2 \, ({\rm H} - {\rm F}) \, g \, (\omega_{\rm H}, \omega_{\rm F}, \tau) \,, \end{split} \tag{2 a}$$

$$R_{\rm HF} = (1/2) \gamma_{\rm H}^2 \Delta M_2 ({\rm H} - {\rm F}) g'(\omega_{\rm H}, \omega_{\rm F}, \tau),$$
 (2b)

$$\begin{split} R_{\mathrm{FF}} &= (1/2) \gamma_{\mathrm{F}}^2 \Delta M_2 (\mathrm{F} - \mathrm{H}) \, g \, (\omega_{\mathrm{F}}, \omega_{\mathrm{H}}, \tau), \\ &\approx (4/3) (1/2) \gamma_{\mathrm{H}}^2 \, \Delta M_2 \, (\mathrm{H} - \mathrm{F}) \, g \, (\omega_{\mathrm{H}}, \omega_{\mathrm{F}}, \tau), \ (2\,\mathrm{c}) \end{split}$$

$$\begin{split} R_{\rm FH} &= (1/2) \, \gamma_{\rm F}^2 \, \Delta M_2 \, ({\rm F-H}) \, g' \, (\omega_{\rm F}, \omega_{\rm H}, \tau) \\ &\approx (4/3) \, R_{\rm HF} \, . \end{split} \tag{2 d}$$

Here.

$$\begin{split} f\left(\omega_{I},\tau\right) &= \tau/(1+\omega_{I}^{2}\,\tau^{2}) + 4\,\tau/(1+4\,\omega_{I}^{2}\,\tau^{2}), \quad (3\,\mathrm{a}) \\ g\left(\omega_{I},\omega_{J},\tau\right) &= \tau/\left\{1+\left(\omega_{I}-\omega_{J}\right)^{2}\,\tau^{2}\right\} + 3\,\tau/\left(1+\omega_{I}^{2}\,\tau^{2}\right) \\ &+ 6\,\tau/\left\{1+\left(\omega_{I}+\omega_{J}\right)^{2}\,\tau^{2}\right\}, \quad (3\,\mathrm{b}) \end{split}$$

$$g'(\omega_I, \omega_J, \tau) = -\tau / \{1 + (\omega_I - \omega_J)^2 \tau^2\} + 6\tau / \{1 + (\omega_I + \omega_J)^2 \tau^2\}.$$
 (3c)

 $\Delta M_2(I-J)$ represents a reduction of the second moment of I spins caused by averaging out the dipolar interactions between I and J spins through the $\mathrm{NH_4^+}$ ion reorientations. γ_I and ω_I are the gyromagnetic ratio and Larmor frequency of I spins, respectively. τ stands for the correlation time of the

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 NH_4^+ ion motion and the Arrhenius relation is usually assumed for τ :

$$\tau = \tau_0 \exp(E_a/RT). \tag{4}$$

By adjusting ΔM_2 (H – H), ΔM_2 (H – F), τ_0 and E_a to fit the calculated T_1 to the observed one, the motional parameters of the NH₄⁺ ions were determined: ΔM_2 (H – H) = 46 × 10⁻⁸ T², ΔM_2 (H – F) = 2.4 × 10⁻⁸ T², E_a = 22 kJmol⁻¹, and τ_0 = 2.0 × 10⁻¹⁴ s. The calculated T_1 curves are shown in Figure 1. The ΔM_2 (H – H) and ΔM_2 (H – F) values from the T_1 analysis are in agreement with 46 × 10⁻⁸ and 4.2 × 10⁻⁸ T², respectively, obtained from the M_2 calculations. Therefore, the relaxation process of both the ¹H and ¹⁹F nuclei in NH₄SnF₃ is assignable to the isotropic reorientation of the cations. The deep T_{1H} minimum is mainly due to the contribution of the homonuclear ¹H – ¹H dipolar interaction, and the T_{1s} minimum at 165 K originates from the heteronuclear ¹H – ¹⁹F dipolar coupling through the $\tau/\{1+(\omega_I-\omega_J)^2\tau^2\}$ term.

From the standpoint of the site symmetry at the NH_4^+ ion site alone it would be expected for the NH_4^+ ions to undergo a reorientation about the C_3 axis in preference to the overall reorientation. However, its possibility is experimentally excluded. This fact indicates that the three hydrogen bonds of the NH_4^+ ion out of the C_3 axis hinder the cation reorientation. In fact, Knop et al. showed from their infrared spectroscopic study that the hydrogen bond on the C_3 axis is much weaker than the remaining three hydrogen bonds [9]. The E_a value of 22 kJmol $^{-1}$ is much smaller than 39 kJmol $^{-1}$ for $\mathrm{NH}_4\mathrm{F}$, in which the NH_4^+ ion forms tetrahedrally four normal hydrogen bonds with an N ... F distance of 271 pm [10], and is much larger than 8.6 kJmol $^{-1}$ for cubic (NH_4)₂SiF $_6$, in which the NH_4^+ ion forms four trifurcated hydrogen bonds with an N ... F distance of 300 pm [7]. In $\mathrm{NH}_4\mathrm{SnF}_3$, the NH_4^+ ion forms three normal and a trifurcated hydrogen bond, and its E_a value is between the two limiting values mentioned above.

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